Note Added in Proof. After this work was completed and in press, a particularly attractive group theoretical treatment of methyl relaxation has appeared (G. B. Matsen, J. Chem. Phys., in press). This treatment leads to spectral density functions identified with symmetry labels rather than with auto- and cross-correlation. The advantage is to provide a physical interpretation for the pre- and postexponential terms in methyl relaxation behavior characterized by multiple exponentials. Subsequent work from this laboratory will use this development.

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# Proton Magnetic Resonance and Conformational Energy Calculations of Repeat Peptides of Tropoelastin: the Tetrapeptide 

Md. Abu Khaled, V. Renugopalakrishnan, and Dan W. Urry*<br>Contribution from the Laboratory of Molecular Biophysics and the Cardiovascular Research and Training Center. University of Alabama Medical Center, Birmingham, Alabama 35294. Received April 8, 1976


#### Abstract

The detailed conformation of a tetrapeptide of tropoelastin, $t$-Boc-L-Val - L- Pro $_{2}$ - $\mathrm{Gly}_{3}$ - $\mathrm{Gly}_{4}$ - OMe in $\mathrm{CDCl}_{3}$, has been obtained from a combined analysis of ${ }^{1} \mathrm{H}$ NMR spectra and conformational energy calculations. The observations of Gly ${ }_{3}$ and Gly 4 methylene protons as ABX spin systems indicate a fixed conformation similar to a cyclic peptide stabilized by hydrogen bond formation. Temperature dependence and solvent perturbation of NH protons and conformational energy calculations each showed the presence of a $\beta$-turn, a ten atom hydrogen-bonded ring involving the $\mathrm{Gly}_{4} \mathrm{NH}$ and $\mathrm{Val}_{1} \mathrm{C}=\mathrm{O}$, and a segment of an antiparallel $\beta$-pleated sheet stabilized by a hydrogen bond between the Val, NH and the Gly $\mathrm{N}_{4} \mathrm{C}=$ O. Conformational angles obtained from the observed ${ }^{3} J_{\alpha \mathrm{CH}} \mathrm{NH}$ coupling constants and from conformational energy calculations were in good agreement. The secondary structure of this tetramer is shown to be the same as previously proposed for the high polymer of the tetramer in water at elevated temperature.


An understanding of the essential functional role of elastin in vascular wall and its molecular pathology in atherosclerosis requires an understanding of its conformation and dynamics.

Structural features at the molecular level are largely responsible for the dynamic properties and interactions of elastin. Gray, Sandberg, and co-workers ${ }^{1.2}$ have shown that soluble
tropoelastin, the precursor protein of fibrous elastin, ${ }^{3-6}$ contains repeating sequences, a tetrapeptide ( $\mathrm{L}-\mathrm{Val}_{1}-\mathrm{L}-\mathrm{PrO}_{2}-$ $\mathrm{Gly}_{3}-\mathrm{Gly}_{4}$, VPGG), a pentapeptide (L-Val $-\mathrm{L}-\mathrm{Pro}_{2}-\mathrm{Gly}_{3}-\mathrm{L}-$ $\mathrm{Val}_{4}-\mathrm{Gly}_{5}$, VPGVG), and a hexapeptide (L-Ala 1 - $\mathrm{L}-\mathrm{Pro}_{2}-$ $\mathrm{Gly}_{3}$-L-Val ${ }_{4}$-Gly ${ }_{5}$-L-Val ${ }_{6}$, APGVGV). The repeat peptides, their oligomers, and high polymers have been synthesized in this laboratory. ${ }^{7-9}$ From extensive studies by NMR using temperature coefficient ${ }^{10-13}$ and solvent perturbation ${ }^{14-16}$ methods, it has been shown that there is a $\beta$-turn in each monomeric unit of these repeat peptides as well as additional secondary structural features and new $\beta$-spiral type conformations for the high polymers of the penta- and hexapeptides of elastin have been proposed. ${ }^{17-19}$ Recently conformational energy calculations on the pentamer ${ }^{20}$ and nuclear Overhauser enhancement measurements on all three repeat peptides ${ }^{21}$ have confirmed that this $\beta$-turn, containing $\mathrm{Pro}_{2}$ and $\mathrm{Gly}_{3}$ at the corners of the ten-membered ring, is type II. Of special interest to the present work, Dale and Titlestad ${ }^{22}$ have reported the ${ }^{1} \mathrm{H}$ NMR spectral behavior of methylene protons ( $\mathrm{CH}_{2}$ ) of glycine or sarcosine contained in cyclic oligopeptides where the $\mathrm{CH}_{2}$ protons appear as ABX or AB spin patterns. The ring constraints make the methylene protons nonequivalent ${ }^{22}$ and result in the $A B$ type spin pattern, Therefore, if the glycine protons in the repeat peptides of elastin are relatively fixed by a $\beta$-turn and additional secondary structural features, they should also be nonequivalent and exhibit $A B X$ spin patterns.

Useful information relative to conformation can be derived by analyzing the ABX spectrum and by using the Karplus relationship ${ }^{23,24}$

$$
\begin{equation*}
{ }^{3} J=A \cos ^{2} \theta+B \cos \theta+C \ldots \tag{1}
\end{equation*}
$$

where $A, B$, and $C$ are constants. Several workers ${ }^{25-29}$ have obtained the values $A, B$, and $C$ for peptide systems by semiempirical and theoretical methods. The widespread application of this relationship has been to evaluate the angles $\phi$ and $\chi$ of peptide backbone and side chains ${ }^{30,31}$ by observing ${ }^{3} J_{\alpha \mathrm{CH}-\mathrm{NH}}$ and ${ }^{3} J_{\alpha \mathrm{CH}-\beta \mathrm{CH}}$, respectively. The limitations of the method are due to the approximate $\cos ^{2} \theta$ dependence of the Karplus relation and the multiple numerical solutions of $\theta$ for a single value of ${ }^{3} J$. Therefore, studies of peptide conformations have effectively utilized a combination of the observed ${ }^{3} J$ values and conformational energy calculations ${ }^{32-36}$ to limit the range of possible conformations. This approach proved to be particularly successful for the valinomycin- $\mathrm{K}^{+}$complex, ${ }^{36}$ where the x -ray structure obtained some years later showed no atom of this dodecamer to be out of place by more than $0.5 \AA .{ }^{37}$ In this study we investigate the detailed conformation of the tetramer of elastin, VPGG, by the combined results of ${ }^{1} \mathrm{H}$ NMR spectral analysis and conformational energy calculations. The conformational energy calculations utilize the "free space approximation" and the partitioned potential energy method. Such calculations have been applied in the past to many biologically important molecules. ${ }^{38-44}$ Recently it has been argued ${ }^{45}$ that conformational energy calculations in vacuo approximate conformations observed in solvents of the low polarity, i.e., of low dielectric constants. In this paper we report the congruity of calculated low-energy conformations and ${ }^{1} \mathrm{H}$ NMR conformational studies in chloroform for a tetrapeptide of elastin. While previous studies on the tetrapeptide of tropoelastin have already led to a proposal for the occurrence of a type II $\beta$-turn involving a $\mathrm{Val}_{1}-\mathrm{C}-\mathrm{O} \cdots \mathrm{HN}-\mathrm{Gly}_{4}$ hydrogen bond and placing $\mathrm{Pro}_{2}$ and $\mathrm{Gly}_{3}$ at the corners ${ }^{7} \cdot 16.21$ and under special circumstances the occurrence of an additional $\mathrm{Val}_{1}$ $\mathrm{NH} \cdots \mathrm{O}-\mathrm{C}-\mathrm{Gly}_{4}$ hydrogen bond, ${ }^{19}$ the present effort confirms these proposals mainly by means of detailed analysis of the glycyl $\mathrm{CH}_{2}$ coupling constants in chloroform and by conformational energy calculations which together provide specification of preferred dihedral angles and a plotted perspective of the conformation.

## Experimental Section

The tetrapeptide of elastin, $t$-Boc- Val $_{1}$ - $\mathrm{Pro}_{2}$-Gly ${ }_{3}$-Gly $y_{4}-\mathrm{OMe}$, was synthesized in this laboratory. ${ }^{1} \mathrm{H}$ NMR spectra were obtained with a 0.1 M solution in $\mathrm{CDCl}_{3}$. To facilitate spectral analysis, $20 \%$ by volume $\mathrm{C}_{6} \mathrm{D}_{6}$ was added (see Figure 1 A and B ). ${ }^{1} \mathrm{H}$ NMR measurements were made on a Varian HR-220 spectrometer operating at a probe temperature of $21^{\circ} \mathrm{C}$ and equipped with an SS- 100 computer system. Simulated spectra were obtained by using a Varian data machine spin simulation program. All the double resonance experiments were performed on a JEOL PS-100 spectrometer operating at a probe temperature of $22^{\circ} \mathrm{C}$ and in the internal lock mode. Variable temperature experiments were made with the PS-100 spectrometer equipped with a JEOL JNM VT- 38 temperature controller.

Conformational Energy Calculations. The partitioned potential energy method has been used for the conformational energy calculations, the program for which originated in the Ramachandran laboratory ${ }^{46}$ and has since been extensively modified. ${ }^{47}$ Total conformational energy, $E_{\text {Total }}$, is expressed as a sum of four separate terms:

$$
\begin{equation*}
E_{\text {Tola) }}=E_{\text {Yan der Waals }}+E_{\text {electroslatic }}+E_{\text {lorsion }}+E_{\mathrm{H} \text {-bonding }} \tag{2}
\end{equation*}
$$

$E_{\text {van der }} W_{\text {aals }}$ was calculated using Buckingham's " 6 -exp" potential function with parameters suggested by Ramachandran and Sasisekharan. ${ }^{38} E_{\text {electrostatic }}$ was calculated assuming a dielectric constant of unity, up to the monopole term, ${ }^{48}$ using ab initio minimal basis set (STO-3G) net charges for glycyl, L-valyl, and L-prolyl residues and $\mathrm{OCH}_{3}$ group, whereas CNDO/ $2^{44}$ net charges were used for the tert-butyloxycarbonyl group. Net charges for glycyl, L-valyl, and L-prolyl residues were derived from an ab initio study of $N$-acetyl-glycine- $N$-methylamide, ${ }^{45} \mathrm{~N}$-acetyl-L-valineamide, and $N$-formyl-L-prolineamide, ${ }^{49}$ Net charges for the $\mathrm{OCH}_{3}$ group were taken from an ab initio study of methanol ${ }^{50}$ and for $t$-Boc groups were derived from a CNDO/2 calculation on $t$-Boc-L-Val $\mathrm{NH}_{2}$. The net charges for VPGG were assembled from those of the fragments, taking care to preserve the overall electroneutrality of the peptide. $E_{\text {torsion }}$ was calculated using a threefold torsional potential with torsional barriers of 0.6 and $0.2 \mathrm{kcal} / \mathrm{mol}$, respectively, for $\mathrm{C}^{\alpha}-\mathrm{N}$ and $\mathrm{C}^{\alpha}-\mathrm{C}^{\prime}$ bonds. ${ }^{51}$ $E_{\mathrm{H} \text {-bonding }}$ was calculated using the following empirical potential function derived by Ramachandran et al. ${ }^{52}$

$$
\begin{equation*}
E_{\mathrm{H} \text {-bonding }}=E_{\text {min }}+p_{1} \Delta^{2}+q_{1} \theta^{2} \tag{3}
\end{equation*}
$$

with $\Delta=R-R_{\text {min }}$ where the best values of the parameters $E_{\text {min }}, p_{1}$, and $q_{1}$ were obtained by comparing the theoretically expected distribution in hydrogen bond length $R$ and hydrogen bond angle, $\theta$, i.e., angle between $\mathrm{N} \cdots \mathrm{O}$ and $\mathrm{N}-\mathrm{H}$ directions, with existing experimental data. The best values obtained ${ }^{52}$ were, $p_{1}=25$ and $q_{1}=0.001$, with $R_{\text {min }}=2.95 \AA$ and $E_{\text {min }}=-4.5 \mathrm{kcal} / \mathrm{mol}$. The above hydrogen bond potential function has been applied to seven-membered ( $\mathrm{C}_{7}$ ) and ten-membered $\left(\mathrm{C}_{10}\right)$ hydrogen-bonded ring systems, i.e., $\gamma$-turns and $\beta$-turns in peptides. ${ }^{53}$ A fully extended conformation of VPGG was constructed using standard dimensions for the peptide group. ${ }^{38}$ so that all the backbone angles $\phi$ and $\psi$ were equal to $180^{\circ}$. ${ }^{54}$ Peptide groups were assumed to be planar and trans with all values of $\omega$ equal to $180^{\circ}$. For the valyl side chain $\mathrm{C}^{\alpha}-\mathrm{C}^{3}, \mathrm{C}^{13}-\mathrm{C}^{\gamma 1}$, and $\mathrm{C}^{3}-\mathrm{C}^{2}{ }^{2}$ bond lengths were assumed to be $1.54 \AA$ and all $\mathrm{C}-\mathrm{H}$ bonds were taken to be 1.10 $\AA$. Bond angles $\mathrm{N}-\mathrm{C}^{\alpha}-\mathrm{C}^{\beta}, \mathrm{C}^{\alpha}-\mathrm{C}^{3}-\mathrm{C}^{\gamma 1}$, and $\mathrm{C}^{1 \gamma}-\mathrm{C}^{\beta 3}-\mathrm{C}^{\gamma 2}$ in the valyl side chain were taken to be 110,112 , and $112^{\circ}$, respectively, ${ }^{55}$ whereas tetrahedral values of $109.47^{\circ}$ were used for all $\mathrm{C}-\mathrm{C}-\mathrm{H}$ and $\mathrm{H}-\mathrm{C}-\mathrm{H}$ bond angles. Initially a trans conformation with $\chi_{1}{ }^{1}=180^{\circ}$ was assumed for the valyl side chain. ${ }^{56}$ We have assumed a $C^{7}$ exo conformation for L -proline with $\chi_{2}{ }^{1}$, i.e., the $\mathrm{N}-\mathrm{C}^{11}-\mathrm{C}^{3}-\mathrm{C}^{-}$torsion angle, equal to about $-17^{\circ}$ as found in $p$-Br-Cbz-Gly-L-Pro-L-LeuGly $(\mathrm{OH}) .{ }^{57}$ Hydrogen atoms were fixed at standard tetrahedral bond angles of $109.47^{\circ}$ and bond lengths of $1.09 \AA$ to the L-prolyl side chain. The $t$-Boc group was assumed to be staggered about the $\mathrm{C}-\mathrm{O}-\mathrm{C}$ plane and was constructed with the dimensions reported by Kashino et al. ${ }^{58}$ for $N$-tert-butyloxycarbonyl- $S$-benzylcysteinylglycine methyl ester. The methoxy group was similarly assumed to be staggered about the $\mathrm{C}-\mathrm{O}-\mathrm{C}$ plane and was constructed with a bond angle for $\mathrm{C}-\mathrm{O}-\mathrm{C}$ equal to $110^{\circ}$ and a tetrahedral geometry for the methyl group. ${ }^{59} \mathrm{~A}$ bond length of $1.42 \AA$ was assumed for the $\mathrm{C}-\mathrm{O}$ bond. The torsion angles necessary for a description of the backbone peptide chain of VPGG are shown in Figure 5 (below) where $\phi_{1}$ and $\psi_{1}$ torsion angles are designated for $\mathrm{Val}_{1}, \phi_{2}$ and $\psi_{2}$ for $\mathrm{Pro}_{2}, \phi_{3}$ and $\psi_{3}$ for Gly $3_{3}$, and $\phi_{4}$ and $\psi_{4}$ for Gly $y_{4}$ residues. The $\phi$ and $\psi$ rotations around the single


Figure 1. $220-\mathrm{MHz}^{\prime} \mathrm{H}$ NMR spectra of $t$ - $\mathrm{Boc}_{-}-\mathrm{Val}_{1}-\mathrm{Pro}_{2}-\mathrm{Gly}_{3}-\mathrm{Gly}_{4}-\mathrm{OMe}\left(\alpha \mathrm{CH}\right.$ region only): (A) in $\mathrm{CDCl}_{3}$; (B) $80 \% \mathrm{CDCl}_{3}$ and $20 \% \mathrm{C}_{6} \mathrm{D}_{6}$ (by volume); and (C) computer simulated spectrum of $\mathrm{Gly}_{3}$ and $\mathrm{Gly}_{4} \mathrm{CH}_{2}$ protons (analyzed as an ABX spin system).
bonds of the backbone, i.e., $\mathrm{N}-\mathrm{C}^{\alpha}$ and $\mathrm{C}^{\alpha}-\mathrm{C}^{\prime}$ in VPGG, were systematically performed starting with the coordinates of the fully extended conformation generated as described above. Due to the rigid conformation adopted by proline, $\phi_{2}$ is fixed near $-60^{\circ}, 56$ the conformational energy of VPGG is a function of seven torsion angles $\phi_{1}$, $\psi_{1}, \psi_{2}, \phi_{3}, \psi_{3}, \phi_{4}$, and $\psi_{4}$, not considering torsion angles necessary to describe the puckering mode of L-prolyl residue. In a later stage of the calculations the side chain torsion angle $\chi_{1}{ }^{1}$ was varied in order to obtain the preferred conformation of L-Val ${ }_{1}$ side chain.

## Results

Proton Magnetic Resonance. The $220-\mathrm{MHz}{ }^{1} \mathrm{H}$ NMR spectra of the $\alpha$-proton region (expanded to $5 \mathrm{~Hz} / \mathrm{cm}$ ) of $t$ -Boc- $\mathrm{Val}_{1}-\mathrm{PrO}_{2}-\mathrm{Gly}_{3}$-Gly ${ }_{4}$-OMe (VPGG) are shown in Figure 1. Figure 1A shows the signals for $\alpha$ protons obtained in $\mathrm{CDCl}_{3}$, the $\alpha \mathrm{CH}$ of $\mathrm{Pro}_{2}$ at 4.43 ppm and $\alpha \mathrm{CH}$ of $\mathrm{Val}_{1}$ at 4.32 ppm were easily assigned from their fine structure and by double resonance experiments. Irradiation of the Boc- $\mathrm{Val}_{1} \mathrm{NH}$, a $9.5-\mathrm{Hz}$ doublet at 5.51 ppm (not shown in the spectrum), collapsed the doublet of a doublet ( 9.5 and 7.5 Hz ) at 4.32 ppm into a doublet showing a $7.5-\mathrm{Hz}$ coupling between the $\alpha \mathrm{CH}$ and the $\beta \mathrm{CH}$ protons of $\mathrm{Val}_{1}$. The Gly ${ }_{3}$ and Gly ${ }_{4}$ methylene protons $\left(\mathrm{CH}_{2}\right)$ appeared more complicated and overlapped. Addition of $20 \% \mathrm{C}_{6} \mathrm{D}_{6}$ (by volume) resulted in all the glycyl methylene proton signals being well-spread and analyzable (see Figures 1 B and 1 C ), The broad triplet of the $\mathrm{Pro}_{2} \alpha \mathrm{CH}$ moved upfield and overlapped with the signal of the Val ${ }_{1} \alpha \mathrm{CH}$. All the chemical shifts and coupling constants are obtained in this solvent mixture $\left(80 \% \mathrm{CDCl}_{3}+20 \% \mathrm{C}_{6} \mathrm{D}_{6}\right)$ and recorded in Table I.

The assignments of the two NH signals for $\mathrm{Gly}_{3}$ and Gly 4 in $\mathrm{CDCl}_{3}$ were made by means of a $\mathrm{CDCl}_{3} \rightarrow \mathrm{Me}_{2} \mathrm{SO}-d_{6}$ solvent titration and by comparison with previously determined temperature dependences and solvent perturbations. ${ }^{7}$ The temperature coefficients ( $\mathrm{d} \delta / \mathrm{d} T$ ) observed in $\mathrm{CDCl}_{3}$ are as follows: Gly NH (a pseudotriplet) at 7.69 ppm shows $\mathrm{d} \delta / \mathrm{d} T$

Table I. ${ }^{1} \mathrm{H}$ NMR Parameters for $t$-Boc-Val - Pro $_{2}{ }_{2}$ Gly $_{3}$-Gly ${ }_{4}$-OMe in $\mathrm{CDCl}_{3}$ with $20 \% \mathrm{C}_{6} \mathrm{D}_{6}(\mathrm{v} / \mathrm{v})$

$a$ Not analyzed. $c$ Overlapped with OMe signal: $s=$ singlet, $\mathrm{d}=$ doublet, $\mathrm{m}=$ multiplet, $\mathrm{dd}=$ doublet of a doublet, $\mathrm{bt}=$ broad triplet, $\mathrm{pt}=$ pseudotriplet. ${ }^{e}$ Values obtained by ABX spin analysis of $\mathrm{CH}_{2}$ proton signals.
$=0.0041 \mathrm{ppm} /{ }^{\circ} \mathrm{C}$; Gly NH (a broad triplet) at 7.52 ppm shows $\mathrm{d} \delta / \mathrm{d} T=0.0091 \mathrm{ppm} /{ }^{\circ} \mathrm{C}$, and $\mathrm{Val}_{1} \mathrm{NH}$ (a doublet, 9.5 $\mathrm{Hz})$ at 5.44 ppm shows $\mathrm{d} \delta / \mathrm{d} T=0.0052 \mathrm{ppm} /{ }^{\circ} \mathrm{C}$. Solvent perturbation on going from chloroform (with $20 \% \mathrm{C}_{6} \mathrm{D}_{6}$ ) to dimethyl sulfoxide at $21^{\circ} \mathrm{C}$ showed that the broad triplet of the Gly NH at 7.52 ppm is the most perturbed one. In previous studies it was found that the Gly ${ }_{4} \mathrm{NH}$ is shielded from the solvent ${ }^{7,15,60}$ and that this shielding is due to H bding to the Val ${ }_{1} \mathrm{C}=\mathrm{O}$, giving a ten-membered ring ( $\beta$-turn)..$^{7.15}$ Since the peptide proton at 7.69 ppm has a low temperature coeffi-

 The $\alpha \mathrm{CH}$ region is shown with the computer simulated spectrum for $\mathrm{Gly}_{3}$ and $\mathrm{Gly}_{4} \mathrm{CH}_{2}$ protons (analyzed as an AB spin system).
cient ${ }^{10-13}$ and is less solvent perturbed, ${ }^{14,15}$ indicative of H bond formation, we therefore assigned this pseudotriplet signal at 7.69 ppm to $\mathrm{Gly}_{4} \mathrm{NH}$. These assignments were confirmed by a $\mathrm{CDCl}_{3} \rightarrow \mathrm{Me}_{2} \mathrm{SO}-d_{6}$ solvent titration on comparison with previous assignments ${ }^{7}$ made in $\mathrm{Me}_{2} \mathrm{SO}-d_{6}$. It can be seen in Figure 1B that the methylene protons of both $\mathrm{Gly}_{3}$ and $\mathrm{Gly}_{4}$ appeared as two $A B X$ spin systems comprising 16 lines in the spectrum. The signals of $\mathrm{Gly}_{3} \mathrm{CH}_{2}$ are designated as ABX , while those of Gly $\mathrm{CH}_{2}$ as $\mathrm{A}^{\prime} \mathrm{B}^{\prime} \mathrm{X}^{\prime}$ (Figure 1B and 1C). The assignments of ABX and $\mathrm{A}^{\prime} \mathrm{B}^{\prime} \mathrm{X}^{\prime}$ were made by the double resonance experiments, such that irradiation of $\mathrm{Gly}_{3} \mathrm{NH}$ at 7.52 ppm collapses the outer eight lines into a four-line $A B$ spectrum. Similarly, by irradiating $\mathrm{Gly}_{4} \mathrm{NH}^{\mathrm{Nat}} 7.69 \mathrm{ppm}$ the inner eight lines collapsed into a four-line $\mathrm{A}^{\prime} \mathrm{B}^{\prime}$ spectrum. This phenomenon can be best noted in Figure 2, where the two ABX spin systems have been transformed into two $A B$ spin systems ( AB for $\mathrm{Gly}_{3}$ and $\mathrm{A}^{\prime} \mathrm{B}^{\prime}$ for $\mathrm{Gly}_{4}$ ) comprised of eight lines upon deuteration of the NH protons. During this experiment it was also observed that the Gly ${ }_{3} \mathrm{NH}$ proton exchanged with deuterium faster than the other two NH protons in the spectrum. This is in agreement with a structure in which the Gly ${ }_{3} \mathrm{NH}$ proton is solvent exposed. The approximate values of chemical shifts and coupling constants were calculated by analyzing the spectrum as an ABX spin system ${ }^{61}$ and then by using the approximate values in a computer simulated spectrum. The correct values of chemical shifts and coupling constants were recorded (Table I) when the computer simulated spectrum exactly matched with the experimentally obtained ${ }^{1} \mathrm{H}$ NMR spectrum (Figure 1B and 1C).

The conformational angle, $\phi$, for $\mathrm{Val}_{1}\left(\phi_{1}\right)$ was calculated by using the Karplus like relation (see eq 1) with the coefficients of Bystrov et al., ${ }^{28}$ i.e.,

$$
\begin{equation*}
{ }^{3} J_{\alpha \mathrm{CH}-\mathrm{NH}}=9.4 \cos ^{2} \theta-1.1 \cos \theta+0.4 \ldots \tag{4}
\end{equation*}
$$

The torsion angles for $\mathrm{Gly}_{3}\left(\phi_{3}\right)$ and $\mathrm{Gly}_{4}\left(\phi_{4}\right)$ were also calculated by using the coefficients of Bystrov et al. ${ }^{28}$ for the

Table II. Backbone Torsion Angles of VPGG Obtained from
${ }^{1} \mathrm{H}$ NMR and Conformational Energy Calculations ${ }^{a}$

$a$ Torsion angles are given in accordance with IUAPC-IUB ${ }^{54}$ convention. $b$ Angles obtained from the Dreiding model of VPGG.
glycyl residue
${ }^{3} J_{\alpha \mathrm{CH}-\mathrm{NH}}+{ }^{3} J_{\alpha \mathrm{CH}^{\prime}-\mathrm{NH}}=-9.8 \cos ^{2} \phi-1.3 \cos \phi+15.0$

Since the $\cos ^{2}$ dependence gives several values of $\phi$ for a single
 II, which are in agreement with angles obtained by conformational energy calculations. Using these $\phi$ torsion angles, the torsion angles $\psi$ were obtained from the Dreiding model of VPGG. These values are also given in Table II.

Low-Energy Conformations. Energy surfaces corresponding to pairs of torsion angles $\phi_{1}-\psi_{1}, \phi_{3}-\psi_{3}$, and $\phi_{4}-\psi_{4}$ for $\mathrm{Val}_{1}$, $\mathrm{Gly}_{3}$, and $\mathrm{Gly}_{4}$, respectively, are shown in Figure 3. The energy surface $\phi_{1}-\psi_{1}$ for $\mathrm{Val}_{1}$ was initially obtained assuming $\left.\chi\right|^{1}$ equal to $180^{\circ}$. As stated above, in the later stages of the calculations the side chain torsion angle $\chi_{1}{ }^{1}=150^{\circ}$ was found to be preferred. A torsional barrier of $2.8 \mathrm{kcal} / \mathrm{mol}$ had been assumed for the side chain rotation. ${ }^{55}$ A plot of energy as a function of $\chi_{1}{ }^{\prime}$ and $\psi_{2}$ for the valyl side chain and $\operatorname{Pro}_{2}$, respectively, are shown in Figure 4 . The $\phi_{1}-\psi_{1}$ energy surface was recomputed using the new value of $\chi_{1}{ }^{1}=150^{\circ}$ and was found not to vary from that shown in Figure 3A. Since a tetrahedral bond angle of $109.47^{\circ}$ was assumed for the bond angles $\tau$, i.e., $\mathrm{C}^{\prime}-\mathrm{C}^{\alpha}-\mathrm{N}$, an attempt was made to relax this criterion using a total minimization scheme originally de-


Figure 3. (A) $\phi_{1}-\psi_{1}$ energy surface in kilocalories/mole relative to the global minimum marked X . Hydrogen-bonded region is indicated by negative contour level. (B) $\phi_{3}-\psi_{3}$ energy surface in kilocalories/mole relative to the global minimum marked X. Hydrogen-bonded region is indicated by negative contour level. Types I and II $\beta$-turns are also marked. (C) $\phi_{4}-\psi_{4}$ energy surface in kilocalories/mole relative to the global minimum marked X. Hydrogen-bonded region is indicated by negative contour level.
scribed by Boyd ${ }^{62}$ and subsequently adapted for peptides. ${ }^{43}$ The stereoscopic drawing of the minimum energy conformation with the stabilizing hydrogen bonds so obtained is shown in Figure 5. In Table II the torsion angles of the minimum energy conformation are compared with torsion angles derived from proton magnetic resonance studies on VPGG. The torsion angles are given in accordance with the IUPAC-IUB convention. ${ }^{54}$

## Discussion

Two secondary structural features of VPGG can be derived from the ${ }^{1} \mathrm{H}$ NMR spectral analysis as well as from conformational energy calculations of this molecule. The nonequivalence of both the $\mathrm{Gly}_{3} \mathrm{CH}_{2}$ and $\mathrm{Gly}_{4} \mathrm{CH}_{2}$ protons indicates that the molecule takes up a fixed conformation, similar to a cyclic peptide ${ }^{63}$ with the formation of two H bonds. A tenmembered $\left(\mathrm{C}_{10}\right) \mathrm{H}$ bond, ${ }^{64,65}$ commonly referred to as a $\beta$ turn, involving $\mathrm{Val}_{1} \mathrm{C}=\mathrm{O}$ and $\mathrm{Gly}_{4} \mathrm{NH}$ are apparent from the


Figure 4. (A) Energy in kilocalories/mole as a function of $\chi_{1}{ }^{\prime}$ for $\mathrm{Val}_{1}$ side chain rotation. (B) Energy in kilocalories/mole as a function of $\psi_{2}$ for $\mathrm{PrO}_{2}$ assuming $\phi_{2}=-60^{\circ}$. Types I and II $\beta$-turn are also marked.


Figure 5. Stereoscopic perspective for the minimum energy conformation of VPGG. $1 \rightarrow 4$ and $4 \rightarrow 1$ hydrogen bonds are also shown.
temperature dependence and solvent perturbation of both Val $_{1}$ NH and Gly ${ }_{4}$ NH protons (see Results section). This conformation also emerges from an analysis of the results of the theoretical calculations presented in Figures 3 and 4. A $\beta$-turn in a linear peptide can occur in two forms, ${ }^{64,65}$ type I and type II, which are related by an approximate $180^{\circ}$ rotation of the peptide moiety between residues two and three of the $\beta$-turn.

From the values of torsion angles (Table II) $\phi_{2}$ and $\psi_{2}$ of $\mathrm{PrO}_{2}$ (Figure 4B) and $\phi_{3}$ and $\psi_{3}$ of Gly ${ }_{3}$ (Figure 3B), a $\beta$-turn of type II can be assigned to VPGG involving $\mathrm{Pro}_{2}$ and $\mathrm{Gly}_{3}$ at the corners as shown in Figure 5. This observation is in agreement with the proposal of Venkatachalam ${ }^{64}$ and Geddes et al. ${ }^{65}$ that a type II $\beta$-turn is formed when the sequence L-X-Gly occurs at the corners.

Experimentally, nuclear Overhauser enhancement (NOE) measurements on repeat tetra-, penta-, and hexapeptides of elastin ${ }^{21}$ and conformational energy calculations on a repeat pentapeptide of elastin ${ }^{20}$ have demonstrated the occurrence of this $\mathrm{PrO}_{2}$-Gly ${ }_{3}$ type II $\beta$-turn. The theoretical calculations on VPGG have shown that the type I $\beta$-turn is less stable by $\sim 2.1 \mathrm{kcal} / \mathrm{mol}$ than the type II $\beta$-turn (see Figure 4 B ). In Figure 3B type I and type II $\beta$-turns are also indicated and the same approximate energy difference is seen between the two forms. However this result is dependent on the dielectric constant chosen for the calculation of the electrostatic energy component. In view of the ${ }^{1} \mathrm{H}$ NMR studies carried out in $\mathrm{CDCl}_{3}$, which has a low dielectric constant, the comparison of the theoretical calculations with experimental studies is expected to be valid as argued in a previous study. ${ }^{45}$ The observation of geminal coupling constants (see Table I) of -17.0 Hz for the $\mathrm{Gly}_{3} \mathrm{CH}_{2}$ and -18.0 Hz for the $\mathrm{Gly}_{4} \mathrm{CH}_{2}$ protons is also indicative ${ }^{66,67}$ of a fixed structure with intramolecular H bonds. It is observed from the theoretical calculations that the $4 \rightarrow \mathrm{I}$ and $\mathrm{l} \rightarrow 4 \mathrm{H}$ bonds are both nonlinear with hydrogen bond lengths of $2.90-3.05 \AA$ and hydrogen bond angle $\theta$ near $30^{\circ}$. Hydrogen bond energies ranging from -2.50 to $-3.40 \mathrm{kcal} / \mathrm{mol}$ were obtained for the $1 \rightarrow 4$ and $4 \rightarrow 1$ hydrogen bonds, depending on the hydrogen bond length $R$ and hydrogen bond angle $\theta$, thereby contributing significantly to the stability of H -bonded conformations. A similar study of the energetics of $\beta$-turn conformations in a tripeptide with alanyl side chains by Chandrasekharan et al, ${ }^{53}$ suggests that the minimum energy conformation usually has a somewhat nonlinear $H$ bond and the observations on VPGG are in line with their conclusions.

The observations of ${ }^{3} J_{\alpha \mathrm{CH}-\mathrm{NH}}$ for $\mathrm{Val}_{1}$, $\mathrm{Gly}_{3}$, and $\mathrm{Gly}_{4}$ give an approximate torsion angle, $\phi$, for these amino acid residues. Since multiple angles can be derived for a single value of ${ }^{3} J_{\alpha \mathrm{CH}-\mathrm{NH}}$, the most probable values of the torsion angle, $\phi$, have been obtained by comparison with the theoretically calculated torsion angles (Table II). Using these $\phi$ angles and taking into account the formation of the two H bonds (mentioned above), a Drieding model of VPGG was constructed from which the corresponding approximate values of the torsion angle, $\psi$, were obtained and also listed in Table II. The theoretically derived results presented in Figure 3 show a cluster of low energy conformations with slight differences in the backbone torsion angles. Torsion angles corresponding to the minimum energy conformation are included in Table II for comparison with the angles obtained from 'H NMR results. From this comparison it can be observed that the torsion angles $\phi_{1}, \psi_{1}, \phi_{4}$, and $\psi_{4}$ indicate the formation of an antiparallel $\beta$-pleated sheet (Figure 5). The ${ }^{2} J_{A^{\prime} B^{\prime}}$ of -18.0 Hz for $\mathrm{Gly}_{4} \mathrm{CH}_{2}$ protons also indicates that the two methylene protons of Gly ${ }_{4}$ are staggered with respect to its adjacent $\mathrm{C}=\mathrm{O}$ group, ${ }^{67}$ giving an extended type conformation ${ }^{63}$ for the Gly 4 residue. This is also supported by measurement of $\phi_{4}$ from its ${ }^{3} J_{\alpha \mathrm{CH}-\mathrm{NH}}$ values (Tables I and II). The observation of $3_{J_{« \mathrm{CH}-\mathrm{NH}}}$ of 9.5 Hz for $\mathrm{Val}_{1}$ strongly suggests a $\beta$-pleated conformation which in combination with $\mathrm{Gly}_{4}$ (through the formation of two H bonds, discussed earlier) gives rise to an antiparallel $\beta$-pleated sheet ${ }^{32}$ for $t$-Boc-VPGG-OMe in solution. The reason for $\mathrm{Val}_{1}$ and $\mathrm{Gly}_{4}$ residues adopting an antiparallel $\beta$-pleated sheet conformation can be understood from a detailed study of the $\phi_{1}-\psi_{1}, \phi_{3}-\psi_{3}$, and $\phi_{4}-\psi_{4}$ energy surfaces presented in Figure 3. The $\phi_{1}-\psi_{1}$ energy surface (Figure 3A)
is somewhat similar to the $\phi-\psi$ plot for an alanyl residue preceding a prolyl residue (see ref 56 and literature cited therein). The $\phi_{1}-\psi_{1}$ energy surface (Figure 3A) shows an allowed low energy region on the upper left side commonly known as $\beta$ pleated sheet region around $\phi_{1}=-70$ to $-180^{\circ}$ and $\psi_{1}=$ $40-170^{\circ}$ where the minimum occurs. This region is preferentially stabilized by the occurrence of hydrogen bonding. In addition a secondary region, $2-4 \mathrm{kcal} / \mathrm{mol}$ higher than the global minimum, occurs around $\phi_{1}=-180$ to $-85^{\circ}$ and $\psi_{1}=$ -50 to $-90^{\circ}$. The $1 \rightarrow 4 \mathrm{H}$ bond ( $\mathrm{C}_{14}$ ) occurs from a combination of the local conformational feature of $\mathrm{Val}_{1}$ residue and the strong electrostatic and van der Waals attractive interaction between $\mathrm{Val}_{1} \mathrm{NH}$ and $\mathrm{Gly}_{4} \mathrm{C}=\mathrm{O}$. The interaction of the Val, side chain with the prolyl side chain and the bulky tertiary butyl group is in large part responsible for this conformational feature.

The observation of a $7.5-\mathrm{Hz}$ (Table I) coupling constant between $\alpha \mathrm{CH}$ and $\beta \mathrm{CH}$ protons of $\mathrm{Val}_{1}$ indicates either a $55 \%$ resident time of its side chain (angle $\chi_{1}{ }^{1}$ ) in trans form ${ }^{68}$ or the angle $\chi_{1}{ }^{1}$ is locked in an off trans conformation. From the conformational energy calculations for the $\mathrm{Val}_{1}$ side chain, $\chi_{1}{ }^{1}$, shown in Figure 4A, it is evident that the appearance of a low energy conformation around $\chi 1^{1}=150^{\circ}$ is in agreement with the observed value of ${ }^{3} J_{\alpha \mathrm{CH}-\beta \mathrm{CH}}$ of 7.5 Hz and a fixed side chain orientation.

## Conclusion

The results of ${ }^{1} \mathrm{H}$ NMR spectral analysis and conformational energy calculations together show a unique antiparallel $\beta$-pleated sheet conformation for $t$-Boc-VPGG-OMe in $\mathrm{CDCl}_{3}$, with a type II $\beta$-turn involving $\mathrm{Pro}_{2}$ and $\mathrm{Gly}_{3}$ at the corners. From a detailed study of temperature dependence of peptide NH protons and carbonyl carbons of poly VPGG (PTP) it was proposed ${ }^{19}$ that this polypeptide (PTP) in water at an elevated temperature adopts a "cross- $\beta$-structure" showing each tetramer unit essentially as an antiparallel $\beta$ pleated sheet. The results reported here for the conformation of tetramer VPGG in $\mathrm{CDCl}_{3}$ lends support to the previous conclusions of the $\beta$-turn and the proposal of a cross- $\beta$-structure (see Figure 7C of ref 19). The spectral features of $\mathrm{Gly}_{3}$ $\mathrm{CH}_{2}$ and $\mathrm{Gly}_{4} \mathrm{CH}_{2}$ protons are of primary importance in detailing the total conformation of this molecule. The two $\mathrm{CH}_{2}$ groups are quite different. In the case of $\mathrm{Gly}_{3}$, one $\alpha \mathrm{CH}$ is cyclohexane-like ${ }^{69}$ ("quasiequatorial") and the other is "quasiaxial" with respect to the adjacent $\mathrm{C}=\mathrm{O}$ group. In the case of $\mathrm{Gly}_{4}$, the $\alpha \mathrm{CH}_{2}$ protons are more or less staggered with respect to the adjacent $\mathrm{C}=\mathrm{O}$ group. ${ }^{67}$ The chemical shift differences between the two methylene protons of both $\mathrm{Gly}_{3}$ and $\mathrm{Gly}_{4}$ (Table I) show the similarity with those of cyclic peptides containing glycine as one of the residues. ${ }^{63,70,71}$ This is largely due to the magnetic anisotropy effect of the adjacent $\mathrm{C}=\mathrm{O}$ group ${ }^{71,72}$ which shows that the proton in the plane of the $\mathrm{C}=\mathrm{O}$ group is shielded more than the one which is out of the plane. This is in good agreement with the previous findings. ${ }^{63,70-73}$ This effect will find increasing application in future work on the investigation of conformation of peptides in solution.
It bears reemphasis that the secondary structure derived in detail in this work for the tetrapeptide in $\mathrm{CDCl}_{3}$ by proton magnetic resonance and in vacuo by conformational energy calculations is the same as previously derived for the high polymer of the tetrapeptide in water by means of proton and carbon-13 magnetic resonance studies which were specifically designed for determining secondary structure (see Figure 7C of ref 19). Thus, while solvent does vary the probability of occurrence of the conformation, the $\beta$-turn has been shown to be the preferred conformation in solvents as diverse as chloroform, methanol, dimethyl sulfoxide, and water. Specifically for the polytetrapeptide in water, the $\beta$-turn (the ten atom
hydrogen-bonded ring) is approximated to occur $70 \%$ of the time below $50^{\circ} \mathrm{C}$ and $90 \%$ above $50^{\circ} \mathrm{C}$ and the 14 atom hy-drogen-bonded ring is approximated to occur $60 \%$ of the time above $50^{\circ} \mathrm{C}$ (see Table IV of ref 19).

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